



THE UNIVERSITY *of* EDINBURGH

Edinburgh Research Explorer

Decay Constants of B and D Mesons from Non-perturbatively Improved Lattice QCD

Citation for published version:

bowler, K, Del Debbio, L, flynn, J, Iacagnina, G, Iles, V, maynard, C & richards, D 2000, 'Decay Constants of B and D Mesons from Non-perturbatively Improved Lattice QCD', *Nuclear physics b*, vol. 619, no. 1-3, pp. 507-537. [https://doi.org/10.1016/S0550-3213\(01\)00511-9](https://doi.org/10.1016/S0550-3213(01)00511-9)

Digital Object Identifier (DOI):

[10.1016/S0550-3213\(01\)00511-9](https://doi.org/10.1016/S0550-3213(01)00511-9)

Link:

[Link to publication record in Edinburgh Research Explorer](#)

Published In:

Nuclear physics b

General rights

Copyright for the publications made accessible via the Edinburgh Research Explorer is retained by the author(s) and / or other copyright owners and it is a condition of accessing these publications that users recognise and abide by the legal requirements associated with these rights.

Take down policy

The University of Edinburgh has made every reasonable effort to ensure that Edinburgh Research Explorer content complies with UK legislation. If you believe that the public display of this file breaches copyright please contact openaccess@ed.ac.uk providing details, and we will remove access to the work immediately and investigate your claim.



Edinburgh 2000/14
 IFUP-TH/2000-17
 JLAB-THY-00-25
 SHEP 00 08

Decay Constants of B and D Mesons from Non-perturbatively Improved Lattice QCD

UKQCD Collaboration

K.C. Bowler ^a, L. Del Debbio ^{b,1}, J.M. Flynn ^b,
 G.N. Lacagnina ^a, V.I. Lesk ^{b,2}, C.M. Maynard ^a,
 D.G. Richards ^{a,c,d}

^a*Department of Physics & Astronomy, University of Edinburgh, Edinburgh EH9
 3JZ, Scotland, UK*

^b*Department of Physics & Astronomy, University of Southampton, Southampton,
 SO17 1BJ, UK*

^c*Jefferson Laboratory, MS 12H2, 12000 Jefferson Avenue, Newport News, VA
 23606, USA.*

^d*Department of Physics, Old Dominion University, Norfolk, VA 23529, USA.*

Abstract

The decay constants of B, D and K mesons are computed in quenched lattice QCD at two different values of the coupling. The action and operators are $\mathcal{O}(a)$ improved with non-perturbative coefficients. The results are $f_B = 195(6)^{+24}_{-23}$ MeV, $f_D = 206(4)^{+17}_{-10}$ MeV, $f_{B_s} = 220(6)^{+23}_{-28}$ MeV, $f_{D_s} = 229(3)^{+23}_{-12}$ MeV and $f_K = 150(3)^{+12}_{-8}$ MeV. Systematic errors are discussed in detail. Results for vector decay constants, flavour symmetry breaking ratios of decay constants, the pseudoscalar-vector mass splitting and D meson masses are also presented.

PACS numbers:12.38Gc, 13.20Fc, 13.20Hc, 14.40Lb, 14.40Nd

¹ Dipartimento di Fisica, Università di Pisa and INFN Sezione di Pisa, Italy.

² Center for Computational Physics, Tsukuba Daigaku, Tsukuba Ibaraki, 305-8577, Japan

1 Introduction

The accurate determination of the B and D meson decay constants is of profound importance in phenomenology. The combination $f_B\sqrt{B_B}$, for both B_d and B_s mesons, plays a crucial role in the extraction from experimental data of CKM quark mixing and CP violation parameters. The phenomenological parameter B_B describes B^0 – \bar{B}^0 mixing, and is expected to be close to unity. D meson decay constants are needed for calculations based on factorisation of non-leptonic B meson decays to charmed mesons (see [1] for a review and [2] for a recent application).

Numerical simulations of lattice QCD provide a method for computing the requisite matrix elements from first principles. A prime concern in such calculations is the control of discretisation errors, most notably those associated with heavy quarks, given that in practice the b -quark mass is large in lattice units: $am_Q > 1$. One approach, pioneered by Bernard *et al.* [3] and by Gavela *et al.* [4] and used in the present study, is to work with heavy quarks around the charm quark mass and extrapolate to the b mass scale guided by continuum heavy-quark effective theory (HQET). Other techniques use some form of effective field theory directly to reduce the cut-off effects associated with heavy quarks. Examples include lattice HQET [5], non-relativistic QCD (NRQCD) [6] and a re-interpretation of the Wilson fermion action as an effective theory for heavy quarks, developed by El Khadra *et al.* [7,8], known as the Fermilab formalism. A critical discussion of these various methods, their associated systematic errors and a survey of recent results for f_B , is given in a comprehensive review by Bernard [9].

This paper presents decay constants and masses of heavy-light mesons calculated in the quenched approximation to QCD at two values of the lattice coupling, $\beta = 6.2$ and $\beta = 6.0$. The calculation uses a non-perturbatively improved relativistic Sheikholeslami-Wohlert (SW) [10] fermion action and current operators so that the leading discretisation errors in lattice matrix elements appear at $\mathcal{O}(a^2)$ rather than $\mathcal{O}(a)$. However, non-perturbative improvement does not necessarily reduce lattice artefacts at a given β . Given results at two β values, heavy quark mass-dependent discretisation errors can be estimated by combining a continuum extrapolation with the fit in heavy quark mass, although a simple continuum extrapolation of the final result is not attempted.

Details of the lattice calculation and the extraction of decay constants and masses from Euclidean Green functions are described in Section 2. The extrapolations to physical quark masses, both heavy and light, and the heavy quark symmetry (HQS) relationship between pseudoscalar and vector decay constants are discussed in Section 3. Results for the decay constants are presented in Section 4 and summarised here.

$$\begin{aligned}
f_B &= 195(6)_{-23}^{+24} \text{ MeV} & f_{B^*} &= 28(1)_{-4}^{+3} \\
f_D &= 206(4)_{-10}^{+17} \text{ MeV} & f_{D^*} &= 8.6(3)_{-9}^{+5} \\
f_{B_s} &= 220(6)_{-28}^{+23} \text{ MeV} & f_{B_s^*} &= 25(1)_{-3}^{+2} \\
f_{D_s} &= 229(3)_{-12}^{+23} \text{ MeV} & f_{D_s^*} &= 8.3(2)_{-6}^{+5} \\
\\
f_{B_s}/f_B &= 1.13(1)_{-5}^{+1} & f_{B^*}/f_{B_s^*} &= 1.10(2)_{-6}^{+2} \\
f_{D_s}/f_D &= 1.11(1)_{-3}^{+1} & f_{D^*}/f_{D_s^*} &= 1.04(1)_{-4}^{+2} \\
f_K &= 150(3)_{-8}^{+12} \text{ MeV}
\end{aligned}$$

The first error quoted is statistical, the second systematic. Systematic errors are discussed in Section 4, with the main contributions itemized in Table 11.

2 Details of the calculation

2.1 Improved action and operators

In the Wilson formulation of lattice QCD, the fermionic part of the action has lattice artifacts of $\mathcal{O}(a)$ (where a is the lattice spacing), while the gauge action differs from the continuum Yang-Mills action by terms of $\mathcal{O}(a^2)$. To leading order in a the Symanzik improvement program for on-shell quantities involves adding the SW term to the fermionic Wilson action,

$$S_{\text{SW}} = S_W - c_{\text{SW}} \frac{i\kappa}{2} \sum_x \bar{\psi}(x) i\sigma_{\mu\nu} F_{\mu\nu}(x) \psi(x) \quad (1)$$

Full $\mathcal{O}(a)$ improvement of on-shell matrix elements also requires that the currents are suitably improved. The improved vector and axial currents are

$$\begin{aligned}
V_\mu^I(x) &= V_\mu(x) + ac_V \tilde{\partial}_\nu T_{\mu\nu}(x) \\
A_\mu^I(x) &= A_\mu(x) + ac_A \tilde{\partial}_\mu P(x)
\end{aligned} \quad (2)$$

where

$$\begin{aligned}
V_\mu(x) &= \bar{\psi}(x) \gamma_\mu \psi(x) \\
A_\mu(x) &= \bar{\psi}(x) \gamma_\mu \gamma_5 \psi(x) \\
P(x) &= \bar{\psi}(x) \gamma_5 \psi(x) \\
T_{\mu\nu}(x) &= \bar{\psi}(x) i\sigma_{\mu\nu} \psi(x)
\end{aligned}$$

and $\tilde{\partial}_\mu$ is the symmetric lattice derivative. The generic current renormalisation is as follows ($J = A, V$):

$$J^R = Z_J(1 + b_J am_q)J^I \quad (3)$$

where Z_J is calculated in a mass-independent renormalisation scheme.

The bare quark mass, am_q , is

$$am_q = \frac{1}{2} \left(\frac{1}{\kappa} - \frac{1}{\kappa_{\text{crit}}} \right) \quad (4)$$

where κ is the hopping parameter. For non-degenerate currents, an effective quark mass is used in the definition of the renormalised current, corresponding to

$$\frac{1}{\kappa_{\text{eff}}} = \frac{1}{2} \left(\frac{1}{\kappa_1} + \frac{1}{\kappa_2} \right) \quad (5)$$

In this renormalisation scheme, the improved quark mass, as used in chiral extrapolations, is defined as

$$\widetilde{m}_q = m_q(1 + b_m am_q) \quad (6)$$

2.2 The static limit

The static quark propagator is calculated using the method of Eichten [5], keeping only the leading term in the expansion of the propagator in inverse powers of the quark mass:

$$S_Q(\vec{x}, t; \vec{0}, 0) = \left(\theta(t) e^{-m_Q t} \frac{1 + \gamma^4}{2} + \theta(-t) e^{m_Q t} \frac{1 - \gamma^4}{2} \right) \delta^{(3)}(\vec{x}) \mathcal{T}_0(t, 0) \quad (7)$$

where $\mathcal{T}_0(t, 0)$ is the product of temporal links from $(\vec{0}, t)$ to the origin.

$$\mathcal{T}_0(t, 0) = \prod_{\tau=1}^t U_4^\dagger(\vec{0}, t - \tau) \quad (8)$$

The prescription for renormalizing and improving the axial static-light current is largely analogous to the propagating heavy-light current [11] case. The difference is that the temporal static axial current requires covariant derivatives:

$$(A_I^{\text{stat}})_0 = A_0^{\text{stat}} + ac_A^{\text{stat}} \bar{\psi}_l \gamma_k \gamma_5 \frac{1}{2} \left(\overleftarrow{D}_k + \overleftarrow{D}_k^* \right) \psi_Q \quad (9)$$

so that

$$(A_{\text{R}}^{\text{stat}})_0 = Z_A^{\text{stat}}(1 + b_A^{\text{stat}} am_{\text{q}})(A_{\text{I}}^{\text{stat}})_0 \quad (10)$$

The c_A^{stat} term is not implemented in this calculation, but its value is negative which will be significant when the static point is compared to the infinite mass extrapolation.

2.3 Definitions of mesonic decay constants

The pseudoscalar and vector meson decay constants, f_P and f_V , are defined by

$$\langle 0 | A_{\mu}^{\text{R}}(0) | P \rangle = i p_{\mu} f_P, \quad (11)$$

$$\langle 0 | V_{\mu}^{\text{R}}(0) | V \rangle = i \epsilon_{\mu} \frac{M_V^2}{f_V} \quad (12)$$

where $|P\rangle$ is a pseudoscalar meson state with momentum p_{μ} , while $|V\rangle$ is a vector meson, with mass aM_V and polarisation vector ϵ_{μ} . A_{μ}^{R} (V_{μ}^{R}) denotes the renormalised axial (vector) current, here taken to be the renormalised, improved local lattice axial (vector) current, defined via equation (3).

2.4 Simulation details

In this study, gauge field configurations are generated using a combination of the over-relaxed [12,13] and the Cabibbo-Marinari [14] algorithms with periodic boundary conditions at two values of the gauge coupling $\beta = 6/g_0^2$. At each β , heavy quark propagators are computed at four values of the hopping parameter, corresponding to quarks with masses in the region of the charm quark mass. For light quark propagators, three values of κ are used, corresponding to masses around that of the strange quark. Table 1 lists the input and derived parameters. Errors quoted in this and other tables are statistical only unless otherwise specified. The value of the hopping parameter corresponding to zero quark mass, κ_{crit} , is taken from [15]. The strange and normal quark masses are fixed using the pseudoscalar meson masses for the pion and kaon where the lattice spacing has been fixed from the pion decay constant, f_{π} . Statistical errors are estimated using the bootstrap [16] with 1000 re-samplings.

2.5 Improvement coefficients

The improvement program requires values for the current and mass improvement and renormalisation coefficients defined in equations (2), (3) and (6) in

Table 1

Input and derived parameters. The lattice spacing is set by f_π .

	$\beta = 6.2$	$\beta = 6.0$
Volume	$24^3 \times 48$	$16^3 \times 48$
c_{SW}	1.614	1.769
N_{configs}	216	305
a^{-1} (GeV)	2.66^{+7}_{-7}	1.91^{+6}_{-6}
Heavy κ	0.1200, 0.1233, 0.1266, 0.1299	0.1123, 0.1173, 0.1223, 0.1273
Light κ	0.1346, 0.1351, 0.1353	0.13344, 0.13417, 0.13455
κ_{crit}	0.13581^{+2}_{-1}	0.13525^{+2}_{-1}
κ_{n}	0.13577^{+2}_{-1}	0.13519^{+2}_{-1}
κ_{s}	0.13479^{+6}_{-6}	0.1338^{+1}_{-1}

Section 2.1. One would like to use non-perturbative determinations of the coefficients in order to remove all $\mathcal{O}(a)$ errors. However, different non-perturbative determinations may give mixing coefficients differing by terms of $\mathcal{O}(a)$ and normalisation coefficients differing at $\mathcal{O}(a^2)$. Using a consistently determined set of coefficients (that is, applying the same improvement condition at each value of the coupling and a consistent set of improvement conditions for all coefficients) should enable a smooth continuum extrapolation.

Alternatively, perturbation theory may be used [25], although this leaves residual discretisation errors of $\mathcal{O}(\alpha_s^2 a)$. Lattice perturbation theory is improved by using a boosted coupling $g^2 = g_0^2/u_0^4$ [26]. The mean link, u_0 , is related to the plaquette expectation value, $u_0^4 = \langle \text{Re Tr } U_P \rangle / 3$.

Non-perturbative determinations of the improvement coefficients are available from two groups. The ALPHA collaboration have determined the value of c_{SW} [17,18] using chiral symmetry and Ward identities in the Schrödinger Functional (SF) formalism. They have determined c_A [18] and Z_A , Z_V and b_V [19] in the same scheme and have a preliminary determination of c_V [20,21]. Bhattacharya *et al.* [22–24] have determined all the improvement coefficients needed to improve and renormalise quark bilinears, also using Ward identities, but on a periodic lattice with standard sources. They also use the ALPHA value of c_{SW} to improve the action.

The two non-perturbative determinations of the improvement coefficients give very similar values for the renormalisation coefficients (Z 's) and the quark mass constant b_V (Table 2). However, the mixing coefficients c_A and c_V differ greatly for the two values of the coupling used here. This is shown in Figure 1.

Table 2

Improvement coefficients from various determinations. LANL refers to Bhattacharya *et al.* and BPT is boosted perturbation theory.

	$\beta = 6.2$			$\beta = 6.0$		
	LANL	ALPHA	BPT	LANL	ALPHA	BPT
Z_A	0.818(2)(5)	0.807(8)(2)	0.8163	0.807(2)(8)	0.7906(94)	0.8038
b_A	1.32(3)(4)	—	1.24	1.28(3)(4)	—	1.26
c_A	-0.032(3)(6)	-0.038(4)	-0.012	-0.037(4)(8)	-0.083(5)	-0.013
Z_V	0.7874(4)	0.7922(4)(9)	0.7959	0.770(1)	0.7809(6)	0.7820
b_V	1.42(1)(1)	1.41(2)	1.24	1.52(1)	1.54(2)	1.26
c_V	-0.09(2)(1)	-0.21(7)	-0.026	-0.107(17)(4)	-0.32(7)	-0.028

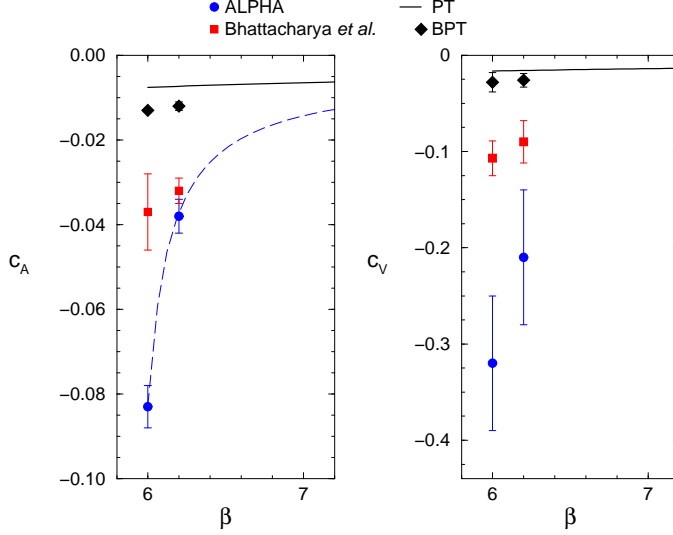


Fig. 1. The improvement coefficients c_A (left) and c_V (right) as functions of the coupling β . Note that the vertical scales in the two plots are different. The values of the coefficients are shown for the two values of the coupling used in this work. The dashed line through the ALPHA c_A points is their interpolating function.

Changes in c_A and c_V can have a particularly large effect on the extracted values of the decay constants when the pseudoscalar and vector meson masses are not small (in lattice units), since the improved current matrix elements are given by,

$$\begin{aligned}
\langle 0|A_4^I|P\rangle &= \langle 0|A_4|P\rangle + c_A \sinh(aM_P) \langle 0|P|P\rangle \\
\langle 0|V_i^I|V, \epsilon\rangle &= \langle 0|V_i|V, \epsilon\rangle + c_V \sinh(aM_V) \langle 0|T_{i4}|V, \epsilon\rangle
\end{aligned} \tag{13}$$

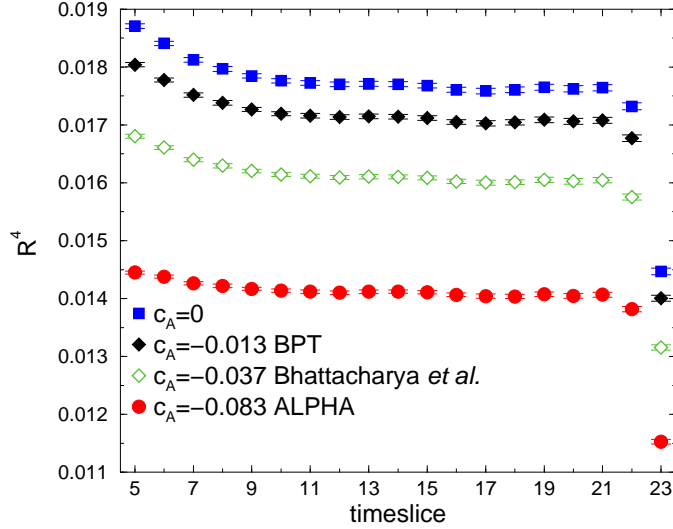


Fig. 2. The mixing of the axial current with the pseudoscalar density. $\beta = 6.0$, $\kappa_H = 0.1123$, $\kappa_L = 0.13344$. The ratio plotted is defined in equation (19).

when the ground state is isolated. In particular at $\beta = 6.0$ the heavy-light meson mass at the heaviest kappa is somewhat bigger than one such that $\sinh aM \sim 1.3$. This is illustrated for the pseudoscalar decay constant in Figure 2. The figure shows the ratio R^4 (between PA and PP correlators, defined in equation (19) and proportional to f_P), for several different values of c_A at $\beta = 6.0$. It can be seen that using the NP value of c_A from the ALPHA collaboration decreases R^4 by $\sim 20\%$ relative to the case $c_A = 0$. The plot also shows the ratio determined using the NP value of c_A from Bhattacharya *et al.*

At $\beta = 6.0$ the disagreement between ALPHA and Bhattacharya *et al.* in the mixing coefficients is striking. Bhattacharya *et al.* try to estimate some of the cut-off effects in their determination of c_A by looking at the difference between two- and three-point derivatives. They also note that since a^2 only halves between $\beta = 6.0$ and $\beta = 6.2$, where disagreement is substantially decreased, even higher order cut-off effects are playing a dominant rôle. Collins and Davies [27] have examined the effect of using higher-order derivatives in determining c_A on the data set used in this work and point out that the difference in the value of c_A from different orders of derivative disappears once the chiral limit is taken. There is clearly a large $\mathcal{O}(a)$ ambiguity in the value of c_A .

Bhattacharya *et al.* have computed all the coefficients needed for this calculation in a consistent fashion and their determinations are used here for that reason. The values of the mixing coefficients c_A and c_V are an order of magnitude larger in the ALPHA determinations than in BPT at $\beta = 6.0$. This has a large impact on the HQS relation between, and scaling of, the decay

constants. When the values of the mixing coefficients from Bhattacharya *et al.* are used the HQS relation is satisfied and there is good scaling.

The ALPHA collaboration compute b_V in the SF scheme at very small quark masses and in Bhattacharya *et al.* the heaviest quark mass is $am_q \sim 0.13$. One may question whether these coefficients can be used around the charm mass: at $\beta = 6.0$, $am_c \sim 0.75$ and at $\beta = 6.2$, $am_c \sim 0.5$. However, the effective normalisation of the vector [28,29] and axial [30] currents has been measured at these quark masses and found to be in good agreement with the determinations of Z_V and b_V ³.

A value for b_m is required to compute the rescaled quark mass used in the chiral extrapolations. Whilst there is a non-perturbative determination [31], it is only at one value of the coupling, and so this work follows [15] and uses the BPT value with which the non-perturbative value is consistent.

Expressions for the static improvement and renormalisation coefficients c_A^{stat} , b_A^{stat} and Z_A^{stat} , have been derived in perturbation theory [11],

$$\begin{aligned} Z_A^{\text{stat}} &= 1.0 + \left(\frac{\ln(a\mu)}{4\pi^2} - 0.137(1) \right) g^2 + \mathcal{O}(g^4) \\ b_A^{\text{stat}} &= \frac{1}{2} - 0.056(7)g_0^2 + \mathcal{O}(g_0^4) \\ c_A^{\text{stat}} &= -\frac{1}{4\pi} \times 1.01(5)g_0^2 \end{aligned} \tag{14}$$

where g_0 is the bare lattice coupling. The one-loop PT and NP renormalisation schemes have been matched at the scale $\mu = m_b$ [32,33] in the $\overline{\text{MS}}$ scheme. The coupling g is then evaluated in the $\overline{\text{MS}}$ scheme [34] at the scale m_b . The values for b_A^{stat} and c_A^{stat} have been evaluated with the boosted coupling at $\beta = 6.0$. The values are

$$\begin{aligned} Z_A^{\text{stat}} &= 0.663_{-15}^{+17} \\ b_A^{\text{stat}} &= 0.4057 \\ c_A^{\text{stat}} &= -0.1354 \end{aligned} \tag{15}$$

where the errors on Z_A^{stat} reflect the uncertainty in $m_b = 4.0 - 4.4$ GeV [34] and in the $\overline{\text{MS}}$ coupling.

³ Indeed Bhattacharya *et al.* compare their results at $\beta = 6.2$ to [28]

Table 3

Pseudoscalar and vector masses in lattice units at $\beta = 6.2$ and $\beta = 6.0$. Fit ranges are 12 – 22 at $\beta = 6.2$ and 10 – 22 (P), 10 – 21 (V) at $\beta = 6.0$.

$\beta = 6.2$				$\beta = 6.0$			
κ_H	κ_L	aM_P	aM_V	κ_H	κ_L	aM_P	aM_V
	0.1346	0.841^{+1}_{-1}	0.871^{+2}_{-2}		0.13344	1.145^{+2}_{-1}	1.188^{+2}_{-2}
0.1200	0.1351	0.823^{+2}_{-1}	0.856^{+2}_{-2}	0.1123	0.13417	1.121^{+2}_{-2}	1.166^{+3}_{-3}
	0.1353	0.817^{+2}_{-1}	0.848^{+3}_{-2}		0.13455	1.110^{+3}_{-2}	1.158^{+4}_{-4}
	0.1346	0.739^{+1}_{-1}	0.775^{+2}_{-2}		0.13344	1.006^{+2}_{-1}	1.056^{+2}_{-2}
0.1233	0.1351	0.721^{+2}_{-1}	0.759^{+2}_{-2}	0.1173	0.13417	0.981^{+2}_{-2}	1.034^{+3}_{-2}
	0.1353	0.714^{+2}_{-1}	0.752^{+3}_{-2}		0.13455	0.969^{+2}_{-2}	1.026^{+4}_{-4}
	0.1346	0.628^{+1}_{-1}	0.673^{+2}_{-2}		0.13344	0.851^{+1}_{-1}	0.915^{+2}_{-2}
0.1266	0.1351	0.609^{+1}_{-1}	0.656^{+2}_{-2}	0.1223	0.13417	0.825^{+2}_{-1}	0.892^{+3}_{-2}
	0.1353	0.602^{+2}_{-1}	0.650^{+3}_{-2}		0.13455	0.811^{+2}_{-2}	0.883^{+4}_{-3}
	0.1346	0.505^{+1}_{-1}	0.563^{+2}_{-2}		0.13344	0.675^{+1}_{-1}	0.759^{+2}_{-2}
0.1299	0.1351	0.484^{+1}_{-1}	0.546^{+2}_{-2}	0.1273	0.13417	0.646^{+2}_{-1}	0.736^{+3}_{-2}
	0.1353	0.476^{+1}_{-1}	0.540^{+2}_{-2}		0.13455	0.631^{+2}_{-1}	0.727^{+4}_{-3}

2.6 Extraction of masses and decay constants of heavy-light mesons

The pseudoscalar meson masses are extracted from the asymptotic behaviour of the two-point correlation functions,

$$C_{PP}^{SS}(t, \vec{p}) = \sum_{\vec{x}} e^{-i\vec{p}\cdot\vec{x}} \langle \Omega_P^S(t, \vec{x}) \Omega_P^{S\dagger}(0, \vec{0}) \rangle \quad (16)$$

$$\xrightarrow{t \rightarrow \infty} \frac{(Z_P^S(\vec{p}))^2}{2aE_P} \cosh(aE_P(t - T/2)) e^{-aE_P T/2}$$

where $T = 48$ and aE_P is the energy of the lowest lying meson destroyed by the operator Ω_P^S and created by $\Omega_P^{S\dagger}$. The superscript S denotes a smeared, or spatially extended, interpolating field operator constructed using the gauge invariant technique described in [35]. $Z_P^S(\vec{p})$ is the overlap of the operator with the pseudoscalar state given by $Z_P^S(\vec{p}) = \langle 0 | \Omega_P^S(0, \vec{0}) | P(\vec{p}) \rangle$. Fit ranges are established by inspection of the time-dependent effective mass. Table 3 shows the best fit masses in each case. A similar procedure is used to extract the masses of vector mesons from correlation functions constructed with the appropriate vector operators.

The decay constants are extracted from the large time behaviour of different two-point correlation functions at zero momentum. The PA correlation function is used for the pseudoscalar decay constant:

$$C_{PA}^{SL}(t, \vec{0}) = \sum_{\vec{x}} \langle A_4^I(t, \vec{x}) \Omega_P^{S\dagger}(0, \vec{0}) \rangle \quad (17)$$

$$\xrightarrow{t \rightarrow \infty} \frac{Z_A^L Z_P^S}{2aM_P} \sinh(aM_P(t - T/2)) e^{-aM_P T/2}$$

where A_4^I is the time component of the improved axial current operator defined in equation (2). The superscript L on the correlator denotes a local operator, in this case the axial current. $Z_A^L = \langle 0 | A_4^I | P(\vec{0}) \rangle$ is the overlap of the local axial operator with the pseudoscalar state, from which the decay constant is extracted using equations (3) and (11).

Extraction of the vector decay constant involves the large time behaviour of the VV correlation function:

$$C_{VV}^{SL}(t, \vec{0}) = \sum_j \sum_{\vec{x}} \langle V_j^I(t, \vec{x}) \Omega_{V_j}^{S\dagger}(0, \vec{0}) \rangle \quad (18)$$

$$\xrightarrow{t \rightarrow \infty} \frac{Z_V^L Z_V^S}{2aM_V} \cosh(aM_V(t - T/2)) e^{-aM_V T/2}$$

where V_j^I is a spatial component of the improved local vector current operator. Again, S denotes a smeared or spatially extended interpolating field operator and L a local operator. The factor Z_V^L is the overlap of the local vector operator with the vector state, $Z_V^L = \sum_r \epsilon_k^r \langle 0 | V_k^I | V(\vec{0}, \vec{\epsilon}) \rangle$, from which the vector decay constant can be extracted via equation (12).

Matrix elements proportional to the decay constants can be extracted from ratios of correlation functions. The ratio

$$R^4 \equiv \frac{C_{PA}^{SL}(t)}{C_{PP}^{SS}(t)} \xrightarrow{t \rightarrow \infty} \frac{Z_A^L}{Z_P^S} \tanh(aM_P(T/2 - t)) \quad (19)$$

is used for the pseudoscalar case and is shown in Figure 2 for various values of c_A . The vector decay constant is determined using

$$\frac{C_{VV}^{SL}(t)}{C_{VV}^{SS}(t)} \xrightarrow{t \rightarrow \infty} \frac{Z_V^L}{Z_V^S} \quad (20)$$

Results for the decay constants are shown in Table 4.

An alternative method is to fit several correlators simultaneously allowing an estimation of the contamination of the ground state signal by excited states.

Table 4

Pseudoscalar and vector decay constants in lattice units. At $\beta = 6.2$ the fit ranges are 14 – 21 (P) and 15 – 23 (V). At $\beta = 6.0$ the fit ranges are 14 – 21 (P) and 16 – 23 (V).

$\beta = 6.2$				$\beta = 6.0$			
κ_H	κ_L	af_P	f_V	κ_H	κ_L	af_P	f_V
	0.1346	0.0892^{+8}_{-8}	9.03^{+10}_{-9}		0.13344	0.1244^{+10}_{-11}	8.79^{+11}_{-9}
0.1200	0.1351	0.0847^{+8}_{-9}	9.28^{+12}_{-12}	0.1123	0.13417	0.1187^{+12}_{-13}	8.94^{+14}_{-12}
	0.1353	0.0832^{+9}_{-9}	9.37^{+14}_{-13}		0.13455	0.1167^{+15}_{-16}	8.89^{+19}_{-14}
	0.1346	0.0873^{+7}_{-8}	7.98^{+8}_{-8}		0.13344	0.1217^{+10}_{-10}	7.80^{+9}_{-8}
0.1233	0.1351	0.0829^{+7}_{-8}	8.17^{+10}_{-9}	0.1173	0.13417	0.1162^{+10}_{-11}	7.89^{+11}_{-10}
	0.1353	0.0814^{+8}_{-9}	8.23^{+11}_{-10}		0.13455	0.1142^{+12}_{-14}	7.84^{+16}_{-11}
	0.1346	0.0847^{+6}_{-7}	6.89^{+6}_{-6}		0.13344	0.1170^{+7}_{-9}	6.74^{+7}_{-6}
0.1266	0.1351	0.0805^{+6}_{-8}	7.00^{+8}_{-7}	0.1223	0.13417	0.1118^{+8}_{-10}	6.79^{+9}_{-8}
	0.1353	0.0790^{+7}_{-9}	7.03^{+9}_{-8}		0.13455	0.1097^{+10}_{-12}	6.72^{+12}_{-9}
	0.1346	0.0804^{+6}_{-7}	5.76^{+5}_{-5}		0.13344	0.1102^{+6}_{-9}	5.62^{+6}_{-5}
0.1299	0.1351	0.0765^{+6}_{-8}	5.79^{+6}_{-5}	0.1273	0.13417	0.1053^{+7}_{-9}	5.59^{+7}_{-6}
	0.1353	0.0751^{+6}_{-9}	5.79^{+7}_{-6}		0.13455	0.1032^{+8}_{-11}	5.51^{+9}_{-7}

An eight-parameter fit is made to C_{PP}^{SS} , C_{PA}^{SL} and C_{PP}^{SL} , allowing for ground and first excited state contributions. Results for the pseudoscalar decay constant are consistent with the single ratio fit. The ratio method is used for central values in the following but the difference from the multi-exponential fits is quoted in Table 14 below as one measure of systematic error.

2.6.1 The static-light axial current

The static-light axial current can be extracted from axial-axial correlation functions. The local-local and local-smeared correlation functions were used in this work, with large-time behaviour given by

$$\begin{aligned}
C_{AA}^{LL}(t, \vec{0}) &= \sum_{\vec{x}} \langle A_4^L(t, \vec{x}) A_4^{L\dagger}(0, \vec{0}) \rangle \xrightarrow{t \rightarrow \infty} (Z_A^L)^2 e^{-a\Delta E t} \\
C_{AA}^{LS}(t, \vec{0}) &= \sum_{\vec{x}} \langle A_4^S(t, \vec{x}) A_4^{L\dagger}(0, \vec{0}) \rangle \xrightarrow{t \rightarrow \infty} Z_A^S Z_A^L e^{-a\Delta E t}
\end{aligned} \tag{21}$$

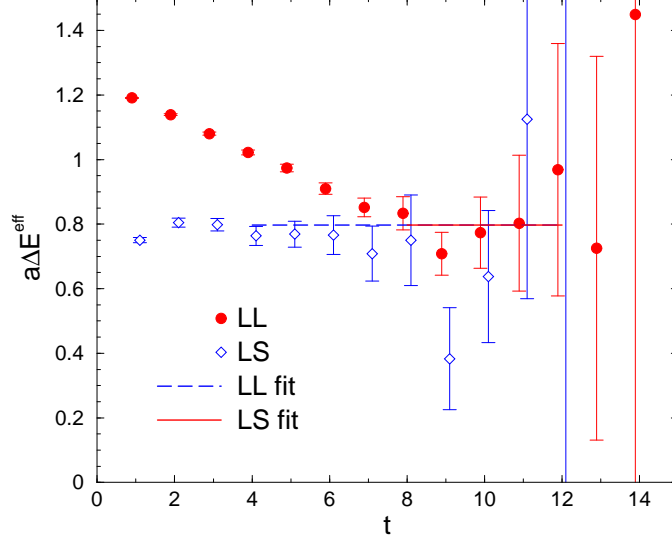


Fig. 3. The effective mass plot for the LL and LS static-light axial correlation function. $\kappa_L = 0.13344$. The lines on the plot show the value of the fit, and the fit ranges.

where ΔE is the unphysical difference between the mass of the meson and the mass of the bare heavy quark. The static-light amplitude, Z_L , is

$$Z_L = af_P^{\text{stat}} \sqrt{\frac{aM_P^{\text{stat}}}{2}} \quad (22)$$

The smearing function is again the one described in [35].

The static-light correlation functions were generated at only one value of the coupling, $\beta = 6.0$, and without the covariant derivative operators necessary to improve the current. The assumption that the matrix element of the improvement term is of the same order of magnitude as the primary term leads to $\mathcal{O}(10\%)$ error in the static point associated with the absence of improvement.

The signal of the static-axial current is very quickly overwhelmed by statistical noise, making the fit difficult. Simultaneously fitting to both the local-local and local-smeared correlator gives a better estimate of the local current. This is shown in Figure 3. The renormalised values of Z_L (without the current improvement term) are shown in Table 5.

2.7 Extraction of masses and decay constants of light-light mesons

The masses and decay constants of the light-light pseudoscalar mesons are extracted from a simultaneous fit to three correlation functions: C_{PP}^{SS} , C_{PP}^{SL} and C_{PA}^{SL} . The smearing function used for the light-lights is the ‘fuzzing’ of

Table 5

Values of $a\Delta E$ and Z_L from simultaneous fits to the LL and LS correlation functions. The fit ranges are 8-12 for the LL and 4-12 for the LS.

κ_L	Z_L	$a\Delta E$
0.13344	0.30(2)	0.80(2)
0.13417	0.28(2)	0.78(3)
0.13455	0.26(3)	0.76(3)

Table 6

The light-light pseudoscalar decay constants. The fit ranges are 8 – 22 at both values of the coupling, except for the lightest mass combination at $\beta = 6.0$ where the fit range was 7 – 23.

$\beta = 6.2$			$\beta = 6.0$		
κ_1	κ_2	af_P	κ_1	κ_2	af_P
0.13460	0.13460	0.0640^{+7}_{-7}	0.13344	0.13344	0.0881^{+10}_{-10}
0.13510	0.13460	0.0610^{+6}_{-7}	0.13417	0.13344	0.0841^{+10}_{-9}
0.13530	0.13460	0.0599^{+6}_{-7}	0.13455	0.13344	0.0821^{+10}_{-10}
0.13510	0.13510	0.0587^{+5}_{-6}	0.13417	0.13417	0.0798^{+9}_{-11}
0.13530	0.13510	0.0575^{+6}_{-6}	0.13455	0.13417	0.0775^{+9}_{-15}
0.13530	0.13530	0.0559^{+6}_{-7}	0.13455	0.13455	0.0730^{+17}_{-18}

reference [36]. The results for the decay constants are shown in Table 6. The results for the pseudoscalar masses are compatible with those in [15] where a similar fit was used.

3 Extrapolation and interpolation in the quark masses

Extrapolation or interpolation in quark masses must be performed to extract physical masses and decay constants. For heavy quarks, the presence of $\mathcal{O}(a^2 m_Q^2)$ lattice artefacts when using the SW action with the NP improved renormalisation scheme imposes a constraint $am_Q < 1$ on the quark masses that can be studied. This limits hadron masses to ~ 2 GeV or slightly greater at $\beta = 6.0$, see Table 10. Input light quark masses are kept above $m_s/2$ to avoid critical slowing down of quark propagator calculations and the possible appearance of finite volume effects. Interpolations to m_c and m_s are needed, together with extrapolations for m_b and the light quarks.

Table 7

The inverse lattice spacing, a^{-1} (GeV)

	$\beta = 6.2$	$\beta = 6.0$
f_π	2.66^{+7}_{-7}	1.91^{+6}_{-6}
r_0	2.91^{+1}_{-1}	2.12^{+1}_{-1}
m_ρ	2.54^{+4}_{-9}	1.89^{+3}_{-5}

3.1 The light-light sector

The dependence of the light pseudoscalar decay constant on the light pseudoscalar meson mass is described by chiral perturbation theory, giving the following ansatz:

$$af_P = c_0 + c_1(aM_P)^2 + c_2(aM_P)^4 \quad (23)$$

The value of af_P for which

$$\frac{af_P}{aM_P} = \left(\frac{f_\pi}{m_\pi} \right)_{\text{expt}} \quad (24)$$

can then be determined and compared to the experimental value of f_π , taken from [34], to fix the lattice spacing. At both values of the coupling, quadratic and linear fits in $(aM_P)^2$ give essentially the same answer for the lattice spacing. Thus the data satisfy lowest-order chiral perturbation theory. Because of quenching and other systematic effects, determinations of the lattice spacing from different quantities disagree. In this work the lattice spacing is set by f_π . Lattice spacings fixed by the Sommer scale, r_0 , [37,38], and m_ρ as determined by [15] are used to estimate systematic error from this source. The values are shown in Table 7.

The light, or ‘normal’, quark mass m_n , defined by $m_n \equiv (m_u + m_d)/2$, and strange quark mass m_s are determined using the lowest-order chiral perturbation theory relation for the mass of a pseudoscalar meson with quark content q_1 and q_2 ,

$$(am_P)^2 = B(a\widetilde{m}_{q_1} + a\widetilde{m}_{q_2}) \quad (25)$$

where the rescaled quark mass \widetilde{m}_q is defined in equation (6). The value of κ_{crit} is taken from a previous UKQCD calculation [15] and is listed in Table 1. The value of the hopping parameter corresponding to the normal quark mass is set

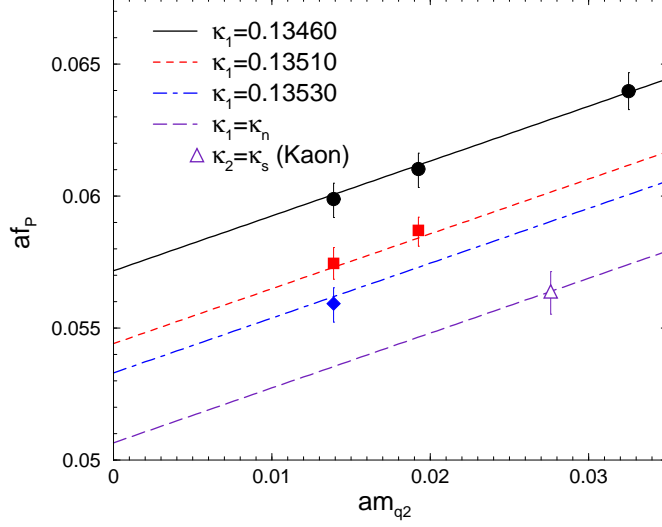


Fig. 4. Dependence of the light pseudoscalar decay constant f_P on quark masses. The lines give f_P as a function of one of the quark masses, $a\tilde{m}_{q_2}$, for fixed values of the other quark mass, $a\tilde{m}_{q_1}$. Filled symbols are the lattice data points. The lowest line shows the extrapolation of $a\tilde{m}_{q_1}$ to the normal quark mass, with the open triangle giving f_K from the position of the kaon on this line.

by the charged pion according to

$$a^2(m_\pi^2)_{\text{expt}} = 2Ba\tilde{m}_n \quad (26)$$

and that for the strange quark mass by

$$a^2(m_K^2)_I = B(a\tilde{m}_n + a\tilde{m}_s) \quad (27)$$

where

$$(m_K^2)_I = \frac{1}{2}(m_{K^\pm}^2 + m_{K^0}^2) \quad (28)$$

with the lattice spacing set by f_π . These hopping parameters are also listed in Table 1.

The ansatz for the dependence of the light pseudoscalar decay constant on quark masses is:

$$af_P = f_0 + f_1(a\tilde{m}_{q_1} + a\tilde{m}_{q_2}) \quad (29)$$

Figure 4 shows the results of fitting to this ansatz and extrapolating to f_K .

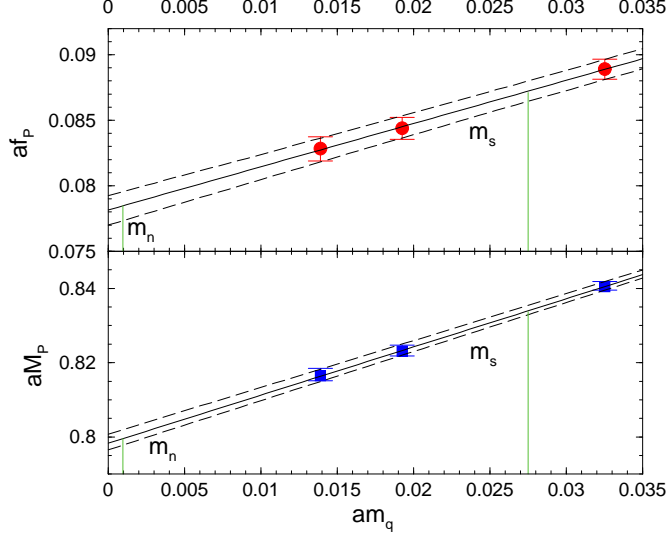


Fig. 5. The chiral extrapolation of the pseudoscalar decay constant (top) and pseudoscalar mass (bottom) against rescaled light quark mass. $\beta = 6.2$ and $\kappa_H = 0.1200$. The vertical lines show the strange and normal quark masses.

3.2 Chiral extrapolations of heavy-light masses and decay constants

A linear dependence of heavy-light masses and decay constants on the light quark mass is assumed:

$$aK_i = \alpha_i + \beta_i a\widetilde{m}_q \quad (30)$$

where K_i is f_P , f_V/a , m_P or m_V . Some example extrapolations are shown in Figure 5. The results for the extrapolated quantities are shown in Tables 8 and 9.

3.3 Heavy quark extrapolations

Heavy Quark Symmetry (HQS) implies asymptotic scaling laws [39] for the decay constants in the infinite heavy quark mass limit. Away from this limit, heavy quark effective theory ideas motivate the following ansätze for the dependence on the heavy meson masses:

$$\Phi_P(M_P) \equiv \Theta(M_B, M_P) f_P \sqrt{M_P} = \gamma_P \left(1 + \frac{\delta_P}{M_P} + \frac{\eta_P}{M_P^2} \right) \quad (31)$$

$$\Phi_V(M_V) \equiv \Theta(M_B, M_V) \frac{M_V}{f_V} \sqrt{M_V} = \gamma_V \left(1 + \frac{\delta_V}{M_V} + \frac{\eta_V}{M_V^2} \right) \quad (32)$$

Table 8

Masses and decay constants at physical light quark masses. $\beta = 6.2$.

κ_H	κ_L	aM_P	af_P	aM_V	f_V
0.1200	κ_n	0.800(2)	0.079(1)	0.832(4)	9.6(2)
	κ_s	0.828(2)	0.0876(8)	0.860(2)	9.1(1)
0.1233	κ_n	0.696(2)	0.077(1)	0.736(3)	8.4(1)
	κ_s	0.726(1)	0.0857(8)	0.764(2)	8.05(9)
0.1266	κ_n	0.583(2)	0.075(1)	0.634(3)	7.1(1)
	κ_s	0.615(1)	0.0832(7)	0.661(2)	6.93(7)
0.1299	κ_n	0.455(2)	0.071(1)	0.523(3)	5.8(1)
	κ_s	0.491(1)	0.0789(7)	0.551(1)	5.77(6)

Table 9

Masses and decay constants at physical light quark masses. $\beta = 6.0$.

κ_H	κ_L	aM_P	af_P	aM_V	f_V
0.1123	κ_n	1.087(4)	0.112(2)	1.138(5)	9.0(2)
	κ_s	1.125(2)	0.122(1)	1.171(3)	8.8(1)
0.1173	κ_n	0.945(3)	0.109(2)	1.005(5)	7.9(2)
	κ_s	0.985(2)	0.119(1)	1.039(3)	7.8(1)
0.1223	κ_n	0.786(3)	0.105(1)	0.862(5)	6.8(1)
	κ_s	0.829(2)	0.115(1)	0.897(3)	6.75(8)
0.1273	κ_n	0.603(2)	0.099(1)	0.704(5)	5.5(1)
	κ_s	0.651(2)	0.1080(9)	0.740(3)	5.60(7)

where Θ denotes logarithmic corrections given at leading order by [40],

$$\Theta(M_B, M) = \left(\frac{\alpha(M)}{\alpha(M_B)} \right)^{2/\beta_0} \quad (33)$$

Here, β_0 is the one-loop QCD beta function coefficient, equal to 11 in the quenched approximation, and $\Lambda_{\overline{\text{MS}}}^{(4)} = 295 \text{ MeV}$ [41].

HQS also relates the pseudoscalar and vector decay constants as follows [40];

$$U(\overline{M}) \equiv \frac{f_V f_P}{\overline{M}} = \left(1 + \frac{8}{3} \frac{\alpha_s(\overline{M})}{4\pi} + \mathcal{O}(1/\overline{M}) \right) \quad (34)$$

where $\overline{M} \equiv (M_P + 3M_V)/4$ is the spin-averaged heavy meson mass. The one-loop factor Θ in equations (31) and (32) cancels in the ratio in equation (34). Higher-order QCD corrections produce the term proportional to α_s . $\tilde{U}(\overline{M})$ is defined to eliminate the radiative corrections in $U(\overline{M})$,

$$\tilde{U}(\overline{M}) \equiv U(\overline{M}) / \left\{ 1 + \frac{8}{3} \frac{\alpha_s(\overline{M})}{4\pi} \right\} \quad (35)$$

Calculated values of $\tilde{U}(\overline{M})$ are fitted to the following parameterisation:

$$\tilde{U}(\overline{M}) = \omega_0 + \frac{\omega_1}{\overline{M}} + \frac{\omega_2}{\overline{M}^2} \quad (36)$$

HQS implies that $\omega_0 = 1$. However, ω_0 can also be left as a free parameter to test the applicability of HQS. Likewise, HQS can be applied to set $\gamma_P = \gamma_V$ in fits using equations (31) and (32). However, higher order QCD corrections would modify this in a similar way to $\tilde{U}(\overline{M})$, that is

$$\frac{\gamma_P}{\gamma_V} = \left(1 + \frac{8}{3} \frac{\alpha_s(\overline{M})}{4\pi} \right) \quad (37)$$

The systematic error arising from extrapolating the decay constants to the b mass is studied in Section 4. The error from truncating the expansion in inverse powers of the heavy mass is considered and the issue of the propagation of discretisation effects under this extrapolation is addressed.

4 Decay constants

The main results for the decay constants are listed in Table 11 and summarised in the introduction. Central values are obtained at $\beta = 6.2$, setting the scale with f_π . These results are discussed in more detail here, starting with f_K .

For f_K the largest source of uncertainty is the choice of quantity used to set the scale. The value of f_K increases by 8% when the scale is set by r_0 and decreases by 4% when it is set by m_ρ . When m_ϕ is used to fix the strange quark mass rather than m_K , the value of f_K decreases by 2.5%. With only two values of the lattice spacing a continuum extrapolation is not attempted. The decay constant is smaller at $\beta = 6.0$ by 2.7%, which is taken as an estimate of its discretisation error. These estimates of systematic errors are combined in quadrature.

The ALPHA collaboration [42] have computed the decay constant, f_K , with the same action as used here, at several values of the coupling. They perform

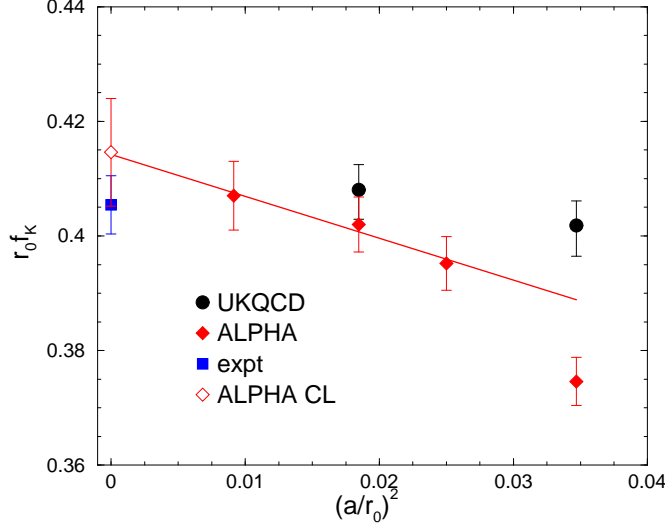


Fig. 6. The dependence of f_K on the lattice spacing.

a continuum extrapolation of $r_0 f_K$ against a^2/r_0^2 . To examine the scaling behaviour of the decay constant in this work, a comparison with the ALPHA results is shown in Figure 6. The line shows the linear extrapolation to the continuum limit (CL) that was performed by ALPHA, excluding the point at the coarsest lattice spacing, $\beta = 6.0$. There are two main differences in the calculations. First, ALPHA use degenerate light quark masses, whereas this work specifically takes into account the non-degeneracy of the quarks. However, ALPHA have checked on the same dataset analysed in this work that using degenerate light quarks has negligible effect. A more obvious difference is in the values of the improvement coefficients, c_A and b_A . As noted earlier, improvement coefficients may differ by $\mathcal{O}(a)$ terms, depending on the improvement condition used. A particular choice of conditions used to determine the coefficients may result, for a particular quantity, in smaller discretisation effects at finite lattice spacing. This seems to be the case for the decay constants when the improvement coefficients of Bhattacharya *et al.* are used rather than those determined by ALPHA, especially at $\beta = 6.0$. However, as the continuum limit is approached the values of f_K from the two studies converge.

The major source of uncertainty in the heavy-light decay constants comes from the determination of the lattice spacing. Additional systematic errors arise from discretisation effects and from the heavy quark extrapolations. Since these latter errors are related to the size of the heavy quark masses, the heavy masses used in this work are listed in Table 10. All these errors will now be discussed.

The ambiguity in the decay constants from the lattice spacing is quite large: there is an overall 16% variation in the value of f_B when the scale is set by m_ρ ,

Table 10

Heavy quark and meson masses used in this work. The table shows the *bare* quark mass in lattice units, the *renormalisation group invariant* quark mass defined by $m_Q^{\text{RI}} = Z_m(1 + b_m am_Q)m_Q$ [43] in GeV and the heavy-light(κ_n) pseudoscalar meson mass in GeV. The scale is set by f_π .

$\beta = 6.2$				$\beta = 6.0$			
κ	am_Q	m_Q^{RI}	M_P	κ	am_Q	m_Q^{RI}	M_P
0.1200	0.485	1.59	2.128(5)	0.1123	0.756	1.26	2.076(8)
0.1233	0.374	1.36	1.851(5)	0.1173	0.566	1.18	1.805(6)
0.1266	0.268	1.06	1.551(5)	0.1223	0.392	0.97	1.501(6)
0.1299	0.168	0.72	1.210(5)	0.1273	0.231	0.65	1.152(4)

Table 11

The decay constants at both β values with the scale set by f_π, r_0 and m_ρ and a quadratic (Q) or linear (L) fit for the extrapolation in inverse heavy meson mass. The quadratic fit uses all four heavy masses, while the linear fit uses the heaviest three. The central values are shown in bold face.

		$\beta = 6.2$				$\beta = 6.0$			
		$f_P(\text{MeV})$		f_V		$f_P(\text{MeV})$		f_V	
		Q	L	Q	L	Q	L	Q	L
f_π	B	195(6)	178(5)	28(1)	29(1)	200(8)	182(7)	26(2)	28(2)
	D	206(4)	207(4)	8.6(3)	8.6(3)	210(4)	211(5)	8.3(3)	8.3(3)
	B _s	220(6)	201(5)	25.2(8)	26.5(8)	222(7)	202(7)	25(1)	26(1)
	D _s	229(3)	230(4)	8.3(2)	8.3(2)	230(4)	230(4)	8.3(2)	8.3(2)
r_0	B	217(5)	200(4)	24.8(8)	25.9(7)	226(6)	207(5)	18.9(8)	20.4(8)
	D	222(3)	223(3)	7.8(2)	7.8(2)	229(3)	230(3)	6.7(2)	6.8(2)
	B _s	240(4)	222(3)	22.9(5)	24.0(4)	247(4)	226(3)	21.9(7)	23.1(5)
	D _s	243(2)	244(2)	7.7(1)	7.7(1)	246(2)	247(2)	7.5(1)	7.5(1)
m_ρ	B	185(4)	168(3)	29.5(9)	31.1(9)	197(6)	179(4)	26(1)	28(1)
	D	199(3)	199(3)	9.1(2)	9.1(2)	208(3)	209(3)	8.4(2)	8.4(2)
	B _s	209(3)	191(2)	26.5(5)	28.0(4)	220(3)	200(2)	24.9(7)	26.5(5)
	D _s	221(2)	222(2)	8.7(1)	8.7(1)	228(2)	228(2)	8.4(1)	8.4(1)

r_0 or f_π (see Table 11). The difference between the decay constants determined at $\beta = 6.2$ and $\beta = 6.0$ is smallest when the scale is set by f_π , as shown in Figure 7. This is not surprising since some systematic uncertainties may cancel

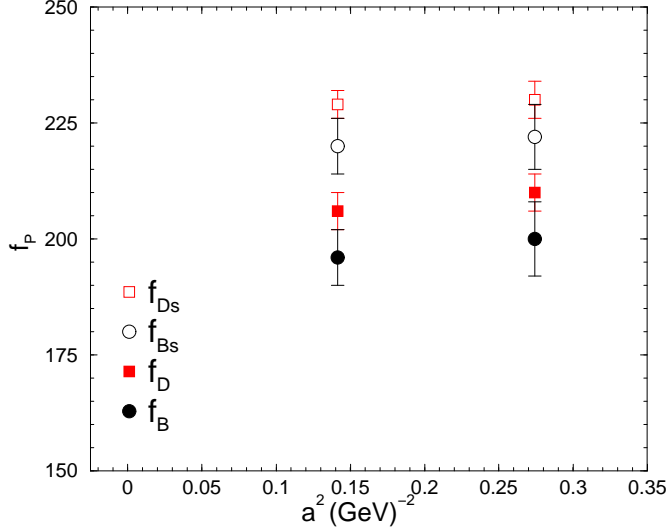


Fig. 7. The dependence of the pseudoscalar decay constants on the lattice spacing. The scale is set by f_π .

in ratios of decay constants. Therefore the central values are quoted using f_π to set the scale.

Discretisation errors are considered next. Although this calculation is $\mathcal{O}(a)$ improved, $\mathcal{O}(a^2)$ effects may still be important at any fixed lattice spacing. In particular, $\mathcal{O}(a^2 m^2)$ effects could well be significant for heavy quarks. With results at only two values of the lattice spacing, a full continuum extrapolation cannot be attempted. The heavy-light decay constant at the finer lattice spacing is taken as the central value, with the result at the coarser spacing used as one indicator of discretisation errors. A fuller investigation of discretisation effects for the heavy quarks follows.

The general characteristic size of discretisation effects in this simulation is investigated by examining the free particle dispersion relation, which is altered in discrete spacetime. The lattice dispersion relation can be written

$$E^2 = M_1^2 + \frac{M_1}{M_2} \vec{p}^2 + \mathcal{O}(p^4) \quad (38)$$

where M_1 is the energy at zero momentum and M_2 is the kinetic mass, defined by $M_2^{-1} = \partial^2 E / \partial p_i^2|_{\vec{p}=0}$. The dispersion relation is investigated by fitting hadronic correlators computed at five different momentum values ($\vec{p}^2 = 0, 1, 2, 3, 4$ in lattice units) using linear and quadratic fits in \vec{p}^2 . The fit for the heaviest quark combination at $\beta = 6.0$ is shown on the left of Figure 8. For lighter mass combinations at $\beta = 6.0$ the signal for highest momentum is very poor, and so in general the fits exclude the $|\vec{p}|^2 = 4$ channel. The right hand side of Figure 8 shows the M_1/M_2 estimates for each heavy quark. Also shown

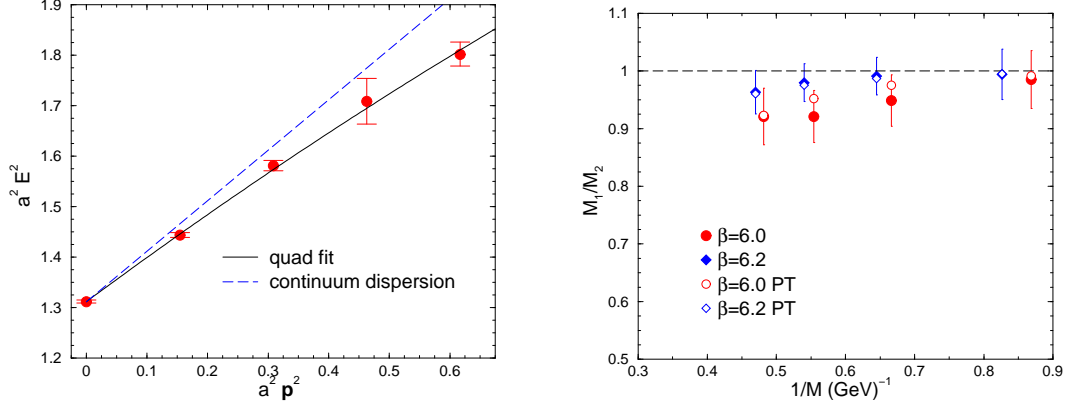


Fig. 8. Left: the pseudoscalar dispersion relation for $\kappa_H = 0.1123$, $\kappa_L = 0.13344$ in lattice units at $\beta = 6.0$. Right: the ratio M_1/M_2 as a function of M_1 .

on the right in Figure 8 is the ratio of M_1/M_2 when M_2 has been determined by the shift in quark mass from m_1 to m_2 [44].

$$M_2^{\text{PT}} = M_1 + (m_2 - m_1) \quad (39)$$

and the quark masses m_1 and m_2 have been calculated from the tree-level relation described in [7,45,46].

$$am_1 = \ln(1 + am_Q) \quad (40)$$

$$\frac{1}{am_2} = \frac{2}{am_Q(2 + am_Q)} + \frac{1}{1 + am_Q} \quad (41)$$

Here M denotes hadron mass and m denotes quark mass. There is good agreement between the tree-level perturbative description of M_2 and the NP determination. At the quark masses used in this work the ratio M_1/M_2 remains close to unity, so that deviations from the continuum dispersion relation are small. Finally, it is worth noting that M_2 is determined at non-zero momentum, which is statistically noisier, and so has statistical errors an order of magnitude larger than for M_1 .

The axial current in this work is normalised according to equations (2,3). As already mentioned the improvement term is proportional to $\sinh(am_P)$ and the mass-dependent normalisation is proportional to am_Q . Bernard [9] proposed the following alternative normalisation, labelled CB,

$$(A_0^R)_{\text{CB}} = Z_A \left[1 + b_A a \bar{m} + 2ac_A \frac{\partial_0 \langle 0|P|P \rangle}{\bar{m} \langle 0|A_0|P \rangle} \right]^{1/2} A_0 \quad (42)$$

where \bar{m} is the average of the bare quark masses in the heavy-light state. The

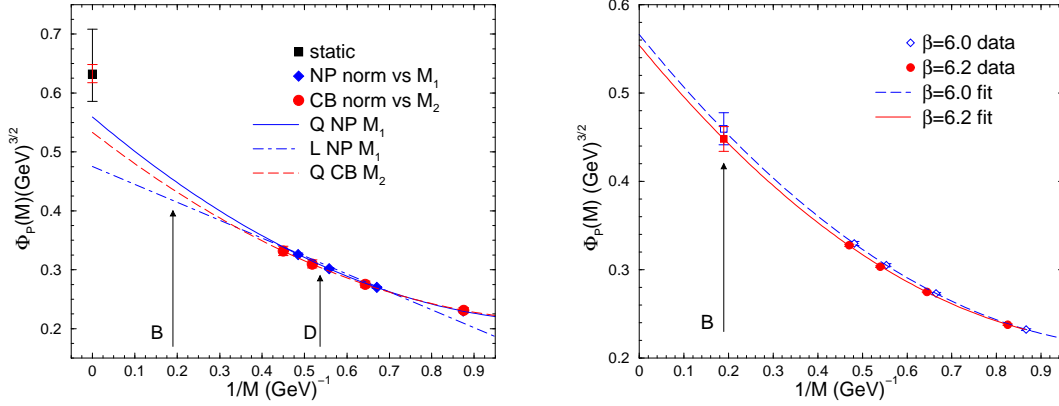


Fig. 9. The left plot shows the extrapolations of the pseudoscalar decay constant in heavy meson mass at $\beta = 6.0$ with different normalisations and masses. The right plot shows a comparison between the extrapolations at both values of the coupling.

CB normalisation differs from the NP improvement scheme normalisation only at $\mathcal{O}(a^2)$. Taking into account an additional normalisation of $\sqrt{4\kappa_Q\kappa_q}$, the CB norm has a finite static limit since $\kappa_Q am_Q \rightarrow 1/2$ in as $\kappa_Q \rightarrow 0$. The CB norm looks very like the normalisation of Kronfeld, Lepage and Mackenzie (KLM norm) [7,26,45,46],

$$\psi \rightarrow \psi' = \psi \sqrt{1 + \mu am} \quad (43)$$

where μ would be the mean link in the tadpole improved tree-level Fermi-lab formalism. In this case μ is a mass-dependent factor that contains the information about improving the current.

Figure 9 shows the mass dependence of the function Φ , defined in equation (31), for a variety of methodologies. The NP norm is shown with both a quadratic fit (Q) to all four masses and a linear (L) fit to the heaviest three. The CB norm data has the extrapolation using the M_2 masses rather than M_1 . It is clear from the figure that the difference between the “NP Q vs M_1 ” and “CB Q vs M_2 ” fits is smaller than the difference between the quadratic and linear fits to NP norm. It should be noted that the effect of changing the normalisation to CB differs from that described in [9] which used the ALPHA determination of c_A and the preliminary NP determination by Bhattacharya *et al.* [22] of b_A .

In the $\mathcal{O}(a)$ improved theory used here, $\mathcal{O}(a^2 m_Q^2)$ lattice artefacts could still be large: the worst case simulated has $a^2 m_Q^2 \sim 0.5$ (see Table 10) and one may worry that mass-dependent discretisation effects, which may not be too large at the charm scale, are enhanced by extrapolation to the bottom scale. Although a Taylor expansion of $(am_Q)^2$ in $1/m_Q$ around the charm quark mass needs large terms to reproduce $(am_b)^2$, the issue is not directly the fate

of $\mathcal{O}(a^2 m_Q^2)$ under extrapolation, but rather the effect of such terms at the charm scale on the extrapolation.

The left hand plot in Figure 9 includes the static point. The smaller set of error bars shows the error from renormalisation, the larger, the statistical uncertainty. The static point is higher than the extrapolations; however, the missing static improvement term would lower the static point, since c_A^{stat} is negative. For this reason an interpolation to the b quark mass using the static and c quark data is not implemented; rather the static point is used as a check of the extrapolation. The result supports the view that, the increase in discretisation error after the heavy quark extrapolation is moderate.

The right-hand side of Figure 9 shows the extrapolation of Φ_P for both values of β in physical units. Although the difference between the curves does indeed grow during the extrapolation, it is smaller than the statistical errors at the bottom scale. The qualitative agreement between the two curves suggests that discretisation errors are not large. To study this further, equation (31) is modified as follows,

$$\Phi(M, a) = \gamma \left(1 + \frac{\delta}{M} + \frac{\eta}{M^2} + \varepsilon(aM)^2 + \zeta(aM)^3 \right) \quad (44)$$

and data from both lattice spacings fitted simultaneously. This uses the two values of a to increase the number of aM values available, allowing the important class of heavy-mass dependent discretisation errors to be studied. The sum of the first three terms in equation (44) has these errors subtracted out and is hereafter referred to as the ‘quasi-continuum’ result. Figure 10 shows the quasi-continuum curve as a function of $1/M$, together with $\Phi(M, a)$ at each β value separately. The curves A and B which blow up as $1/M \rightarrow 0$ are the fits to equation (44) containing lattice artefacts. However, the quasi-continuum result, where these artefacts are subtracted, does not differ greatly from the extrapolations at the individual lattice spacings using equation (31). The quasi-continuum Φ gives $f_B = 186(10)$ MeV, which differs by 5% at $\beta = 6.2$ (the number appearing in Table 14) and 7% at $\beta = 6.0$ from the extrapolations using (31). Excluding the ζ term, or using am_Q instead of aM makes little difference. Using $(aM)^2 \Lambda/M$ instead of $(aM)^2$ also makes no significant difference. Varying all five fit parameters in equation (44) over the region where the chi-squared per degree of freedom increases by up to 1 from its minimum value gives a variation of ± 16 MeV for the quasi-continuum f_B , so the fit is stable.

The systematic error associated with truncating the HQET series is now discussed. First, linear and quadratic fits to the data are compared. The linear fit uses the heaviest three quark masses and drops the η term from equations (31) and (32). This is compared to a quadratic fit using all four masses.

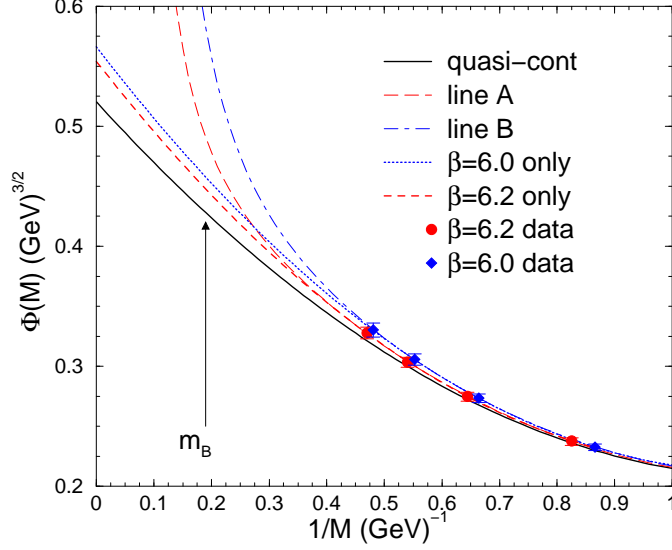


Fig. 10. Effects of lattice artefacts on the heavy quark extrapolation. The solid line shows the quasi-continuum extrapolation described in the text. The dashed and dotted lines show Φ , equation (31), when extrapolating at each value of β separately. Line A(B) shows the function $\Phi(M, a)$, equation (44), at $\beta = 6.2(6.0)$.

The resulting variations, which are shown in Table 11, amount to a relative error of 9% for f_B , indicating that a quadratic term is necessary. A cubic fit is performed to check that a quadratic term is sufficient. There are not enough data points for a cubic fit at each lattice spacing, so the fit combines data from both lattices as before, adding a cubic term ξ/M^3 to $\Phi(M, a)$. This results in a value of f_B which is 5% greater than the quasi-continuum determination. This is also included as a systematic error in Table 14.

The HQS relation between the pseudoscalar and vector decay constants can also be investigated. The quantity $\tilde{U}(\overline{M})$ defined in equation (35) should be equal to unity in the static limit. The extrapolation of $\tilde{U}(\overline{M})$ is shown in Figure 11 and displayed in Table 12. “Q” denotes a quadratic fit to all heavy quarks whereas “L” denotes a linear fit to the heaviest three. The extrapolation for $\beta = 6.0$ displays the expected static limit, albeit with large errors. At $\beta = 6.2$, $\tilde{U}(\overline{M})$ from the quadratic fit deviates from unity in the static limit. On the finer lattice, discretisation errors are smaller and so one might expect better agreement in the static limit. However, $\tilde{U}(\overline{M})$ depends on the ratio of the axial to vector currents and is particularly sensitive to the mixing coefficients, c_A and c_V , which are the most poorly known of the improvement coefficients. Different determinations of c_V vary by an order of magnitude, even at $\beta = 6.2$, where all the other coefficients are in much better agreement. Indeed, varying c_V by the quoted errors of 25% in Bhattacharya *et al.* varies the value of $\tilde{U}(\overline{M}_\infty)$ by around 6%. Taking into account the uncertainty in the value of c_V , the value of $\tilde{U}(\overline{M}_\infty)$ is consistent with unity.

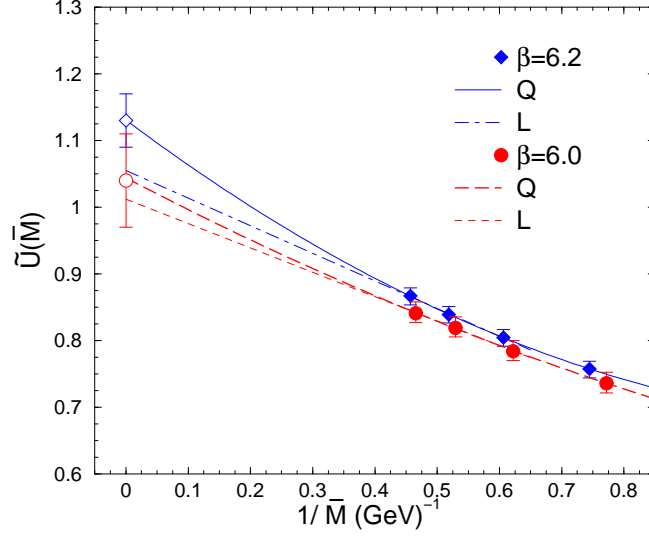


Fig. 11. The quantity $\tilde{U}(\overline{M})$ as a function of the inverse spin averaged mass \overline{M} .

Table 12
 $\tilde{U}(\overline{M})$ as a function of heavy meson mass.

β		\overline{M}_D	\overline{M}_B	\overline{M}_∞
6.2	Q	0.83(1)	1.01(2)	1.13(4)
	L	0.83(1)	0.98(2)	1.05(5)
6.0	Q	0.82(2)	0.96(4)	1.04(7)
	L	0.82(2)	0.94(3)	1.01(4)

Table 13
Flavour-breaking ratios.

β		$\frac{f_{B_s}}{f_B}$	$\frac{f_{D_s}}{f_D}$	$\frac{f_{B^*}}{f_{B_s^*}}$	$\frac{f_{D^*}}{f_{D_s^*}}$
6.2	(i)	1.13(1)	1.11(1)	1.10(2)	1.04(1)
	(ii)	1.11(2)	1.11(1)	1.10(2)	1.05(1)
6.0	(i)	1.11(2)	1.09(1)	1.05(3)	1.00(2)
	(ii)	1.10(2)	1.09(1)	1.05(3)	1.01(2)

The $SU(3)$ flavour breaking ratios are determined by two methods: from the ratio of extrapolated values of the decay constants, labelled (i) in Table 13, or by extrapolating the ratios constructed at each heavy mass, labelled (ii) in the Table. In the ratio of f_{P_s}/f_{P_n} most of the heavy quark mass dependence seems to cancel. The fact that the two methods agree, within statistical errors, suggests that the systematic error in these ratios is small.

Table 14

Percentage systematic uncertainties. Systematic differences are obtained by varying the procedure used to calculate the decay constants. The central values are fixed by using f_π to set the scale at $\beta = 6.2$, performing a quadratic heavy quark extrapolation, taking the central values of the mixing coefficients $c_{A,V}$, fitting correlation functions to a single exponential, and using m_K^2 to set the strange quark mass.

Pseudoscalar	f_B	f_D	f_{B_s}	f_{D_s}	$\frac{f_{B_s}}{f_B}$	$\frac{f_{D_s}}{f_D}$
scale set by r_0	11	8	9	6	-2	-2
scale set by m_ρ	-5	-3	-5	-3	1	1
linear vs quadratic	-9	-	-9	-	-	-
quasi-continuum	-5	-2	-5	-2	-	-
quasi-continuum plus $1/M^3$	5	-	5	-	-	-
$\beta = 6.0$	3	2	1	1	-1	-2
multi-exp	-2	-3	-5	-3	-3	1
strange quark mass from m_ϕ	-	-	-2	-2	-2	-2
coeff. c_A	1	1	1	1	1	-
Vector	f_{B^*}	f_{D^*}	$f_{B_s^*}$	$f_{D_s^*}$	$\frac{f_{B_s^*}}{f_{B^*}}$	$\frac{f_{D_s^*}}{f_{D^*}}$
scale set by r_0	-11	-10	-9	-8	2	2
scale set by m_ρ	6	5	5	5	1	1
$\beta = 6.0$	-7	-4	-2	-1	-5	-4
linear vs quadratic	5	-	5	-	-	-
quasi-continuum	6	3	6	3	-	-
quasi-continuum plus $1/M^3$	-5	-	-5	-	-	-
strange quark mass	-	-	2	2	-3	-2
coeff c_V	3	1	3	1	1	1

The systematic variations of the decay constants are shown in Table 14, in order of importance. The total systematic uncertainty is obtained by combining these in quadrature.

4.1 Comparison with other determinations

A comparison with some other recent quenched results for f_B is shown in Table 15. For a comprehensive review, the reader is referred to the article by Bernard [9]. The results which use the Fermilab formalism [7] are continuum limits. The APE result uses the same action at $\beta = 6.2$ and the UKQCD

Table 15

Comparison with some other recent determinations of f_B . “FNAL” denotes use of the Fermilab formalism [7] and “charm” indicates that the heavy quarks have masses around charm and are then extrapolated to the bottom scale. NB: the World Average is not the average of the numbers displayed below.

	Heavy Quark	f_B (MeV)
El-Khadra <i>et al.</i> [47]	FNAL	$164(8)^{+11}_{-8}$
MILC [48,49]	FNAL	$173(6)^{+16}_{-16}$
JLQCD [50]	FNAL	$173(4)(9)(9)$
UKQCD Tad [51]	charm	$177(17)^{+22}_{-26}$
APE [52]	charm	$179(18)^{+34}_{-9}$
CP-PACS [53]	FNAL	$188(3)(9)$
This Work	charm	$195(6)^{+24}_{-23}$
World Average [9]		$175(20)$

Tad result uses the same gauge configurations as this work but with a tadpole improved value of c_{SW} .

The value of f_B obtained in this work is the highest, although it is compatible with the world average. One reason for this might be the heavy quark method employed. However, other results using the same heavy quark extrapolation have a lower value, while the most recent result using the Fermilab method from the CP-PACS collaboration [53] has a value in agreement within statistical errors. Different heavy quark methods have different associated systematic uncertainties. A detailed analysis of those present in this work has been discussed above, including an estimate of the effects of the heavy quark extrapolation. These uncertainties are also explored in [9], which additionally addresses some of the systematic issues affecting the Fermilab formalism. The value obtained in this work is consistent within these systematic uncertainties with other quenched results.

There have been several recent calculations of f_B with two flavours ($N_F = 2$) of dynamical fermions using NRQCD or the Fermilab formalism. Again, they are reviewed in [9]. The effect of unquenching is found to increase the value of f_B by 10 – 15%.

The only heavy-light decay constant to have been measured experimentally so far is f_{D_s} [54],

$$f_{D_s} = 260 \pm 19 \pm 32 \text{ MeV} \quad (45)$$

Table 16

The hopping parameter corresponding to the charm quark mass. The label FIT refers to the ratio ϵ/λ in equation (46) having a fitted value, whilst BPT refers to the ratio having the BPT value.

	$\beta = 6.2$		$\beta = 6.0$	
a^{-1}	FIT	BPT	FIT	BPT
$a^{-1}(r_0)$	0.12498^{+6}_{-5}	0.12513^{+6}_{-5}	0.11952^{+8}_{-6}	0.12056^{+8}_{-6}
$a^{-1}(m_\rho)$	0.12221^{+7}_{-5}	0.12227^{+7}_{-5}	0.1160^{+1}_{-1}	0.1165^{+2}_{-2}
$a^{-1}(f_\pi)$	0.1232^{+6}_{-5}	0.1233^{+6}_{-6}	0.1164^{+10}_{-11}	0.1171^{+3}_{-10}

With such large uncertainties, the experimental value is consistent with both the unquenched and quenched values of f_{D_s} .

5 Spectroscopic quantities

To determine the value of the hopping parameter, κ_c , corresponding to the charm quark mass requires an interpolation in heavy quark mass. The improved quark mass definition in equation (6) also applies to heavy quarks, so if the meson mass depends linearly on the improved heavy quark mass, the corresponding *bare* quark mass dependence is,

$$aM_H(m_Q) = \rho + \lambda am_Q + \epsilon a^2 m_Q^2 \quad (46)$$

where $\epsilon/\lambda = b_m$. In the NP improved formulation all lattice artefacts of $\mathcal{O}(a)$ have been removed. However, $\mathcal{O}(a^2 m_Q^2)$ effects for the heaviest quarks could be significant. Any contributions at this order would affect the ratio ϵ/λ such that it was no longer equal to b_m . A fit to equation (46) was tried with ϵ/λ fixed to b_m from boosted perturbation theory, and with ϵ/λ allowed to vary freely⁴. The heavy quark dependence can be used to fix κ_c by choosing a particular state (or splitting) to have its physical value. Choosing the pseudoscalar mass to fix κ_c , the spectrum of heavy-light mesons can then be predicted.

For $\beta = 6.2$, the results of using equation (46) are shown in Figure 12. Whilst the value of $\epsilon/\lambda = -0.505(4)$ (labelled FIT in the figure) differs somewhat from the value of $b_m = -0.652$ from BPT, it makes little difference to the value of the κ_c , of order 0.1%. The choice of quantity to set the lattice spacing clearly has a rather large effect. Values for κ_c are shown in Table 16, using the free fit to set the value of the hopping parameter.

⁴ This procedure is entirely equivalent to that of Becirevic *et al.* [52].

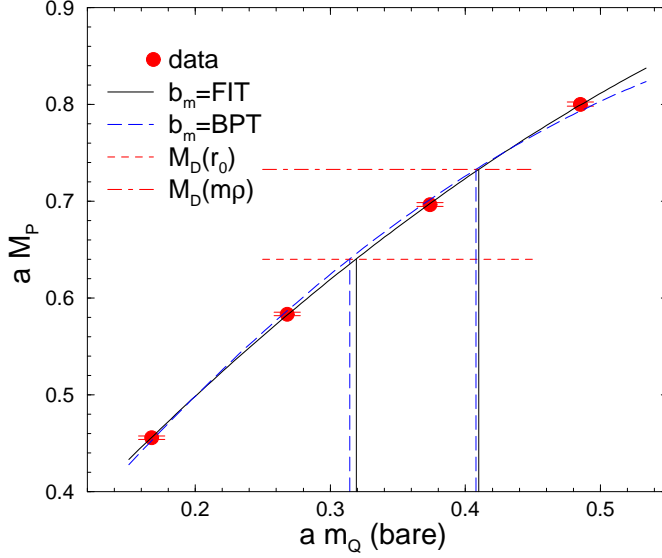


Fig. 12. The pseudoscalar meson vs. bare quark mass for $\beta = 6.2$

For $\beta = 6.0$, the fitted value of $\epsilon/\lambda = -0.384(3)$ differs significantly from the value of $b_m = -0.662$ from BPT. Consequently the fit using the BPT b_m has a very large χ^2 , as the data cannot accommodate a model with such large curvature. However, the difference in κ_c is still small (of order 1%). For $\beta = 6.2$ the heaviest quark has a value of $am_Q = 0.485$ whereas for $\beta = 6.0$ the heaviest quark has $am_Q = 0.775$. Discretisation errors of $\mathcal{O}(a^2m^2)$ could be responsible for modifying the value of the ratio ϵ/λ . The value of κ_c is rather insensitive to ϵ/λ but because the value of am_Q is so large, the improved quark mass changes dramatically. Using the free fit as the preferred method, the value of κ_c is shown in Table 16.

The spectrum of heavy-light charm states is shown in Figure 13 and tabulated in Table 17. With only two values of β a continuum extrapolation is not attempted, but the difference between the two couplings can be used to estimate systematic errors. It is clear from the figure that a large systematic error arises from the definition of the lattice spacing. Depending on which quantity is chosen, different mass splittings appear to be closer to their continuum values. The central values for the spectrum are produced by the following procedures: single exponential fits, using r_0 to set the scale, with ϵ/λ a free parameter in equation (46). The values of the masses at $\beta = 6.2$ are taken to be the central values, with the values at $\beta = 6.0$ used to estimate systematic error. The remaining estimates of systematic error come from using multiple exponential fits, using m_ρ to set the scale and, in the case of states with a strange quark, using m_{K^*} instead of m_K^2 to set the strange quark mass. All these systematic errors are combined in quadrature. The results, with experimental

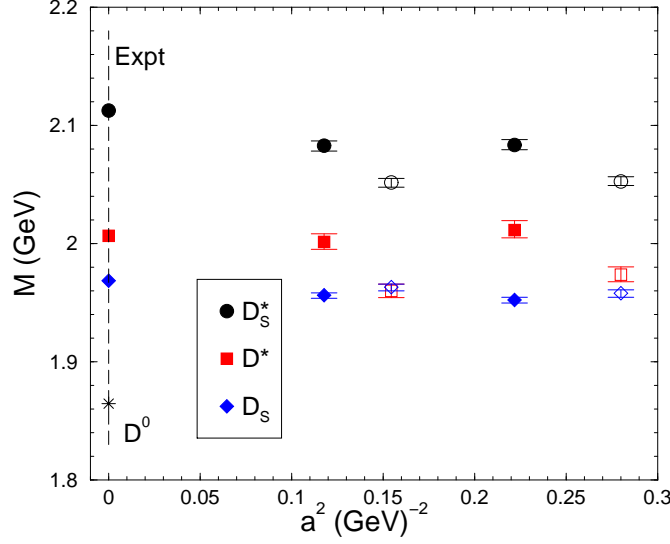


Fig. 13. The scaling of the D meson spectrum. The filled symbols have the lattice spacing set by r_0 , and the open symbols by m_ρ .

comparisons, are:

Lattice, this work	Experiment [34]	
$M_{D_s} = 1.956_{-2}^{+2} {}_{-4}^{+22} \text{ GeV}$	$M_{D_s^+} = 1.9685_{-5}^{+5} \text{ GeV}$	(47)
$M_{D^*} = 2.002_{-6}^{+7} {}_{-41}^{+16} \text{ GeV}$	$M_{D^{*0}} = 2.0067_{-5}^{+5} \text{ GeV}$	
$M_{D_s^*} = 2.082_{-5}^{+4} {}_{-31}^{+22} \text{ GeV}$	$M_{D_s^{*+}} = 2.1124_{-7}^{+7} \text{ GeV}$	

For the lattice results, the first error is statistical and the second systematic. In this calculation the light quark in each meson is a normal quark. Thus the D^* is the isospin-averaged vector state, with mass $M_{D^*} = (M_{D^{*0}} + M_{D^{*\pm}})/2$. With somewhat large systematic errors, the largest of which comes from the lattice spacing, the spectrum is in broad agreement with experiment.

It is also of interest to look at the $M_V - M_P$ splitting as a function of the (pseudoscalar) meson mass. This can be extrapolated in inverse mass to the B scale, and to the infinite quark mass limit where the splitting should vanish. The splittings at both values of β are given in Table 18 and are plotted for $\beta = 6.2$ in Figure 14. There is little difference between linear and quadratic fits, so only the linear fit is shown. The splittings obtained from the linear fit at $\beta = 6.2$, using r_0 to set the scale, are taken as the central values quoted below. The second set of errors is systematic, obtained by combining in quadrature the differences arising from the following variations in procedure (in descending order of importance): using m_ρ rather than r_0 to set the scale; (for the B meson) using a quadratic rather than linear extrapolation; multiple rather

Table 17

The spectrum of heavy-light mesons with different definitions of the lattice spacing in GeV.

state	a^{-1}	$\beta = 6.2$	$\beta = 6.0$
D_s	r_0	1.956^{+2}_{-2}	1.952^{+2}_{-2}
	m_ρ	1.963^{+2}_{-3}	1.958^{+3}_{-3}
	f_π	1.960^{+4}_{-4}	1.958^{+4}_{-3}
D^*	r_0	2.002^{+7}_{-6}	2.011^{+8}_{-7}
	m_ρ	1.960^{+6}_{-6}	1.973^{+7}_{-6}
	f_π	1.973^{+10}_{-10}	1.977^{+14}_{-14}
D_s^*	r_0	2.082^{+4}_{-5}	2.083^{+4}_{-4}
	m_ρ	2.052^{+3}_{-4}	2.053^{+4}_{-3}
	f_π	2.061^{+7}_{-7}	2.056^{+10}_{-10}

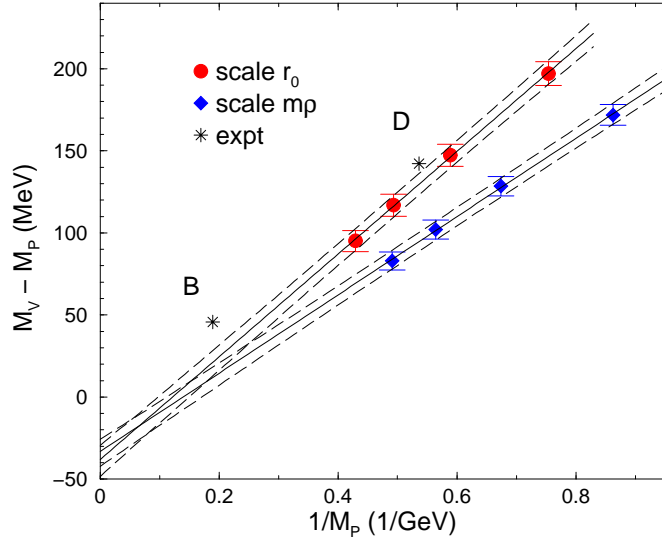


Fig. 14. The hyperfine splitting as a function of M in physical units at $\beta = 6.2$. The circles have the scale set by r_0 and the diamonds by m_ρ .

than single exponential fits; using $\beta = 6.0$ data rather than $\beta = 6.2$ data. The results, with the comparison to experiment, are as follows:

Lattice, this work	Experiment [34]
$M_{D^*} - M_D = 130^{+6+15}_{-6-35}$ MeV	$M_{D^*} - M_D = 142.6 \pm 0.5$ MeV (48)
$M_{B^*} - M_B = 21^{+7+18}_{-8-16}$ MeV	$M_{B^*} - M_B = 45.2 \pm 1.8$ MeV

Table 18

The hyperfine splitting of the heavy-light mesons at the B and D meson scale, in MeV.

state	a^{-1}	$\beta = 6.2$	$\beta = 6.0$
	r_0	130^{+7}_{-6}	140^{+8}_{-6}
$M_{D^*} - M_D$	m_ρ	95^{+6}_{-6}	108^{+7}_{-6}
	f_π	105^{+9}_{-9}	111^{+14}_{-13}
	r_0	21^{+7}_{-8}	32^{+8}_{-8}
$M_{B^*} - M_B$	m_ρ	12^{+6}_{-7}	22^{+8}_{-8}
	f_π	15^{+8}_{-8}	24^{+14}_{-14}

6 Conclusions

The decay constants of heavy-light and light mesons have been determined in the quenched approximation at two values of the coupling. The action and currents have been fully NP $\mathcal{O}(a)$ improved. Good scaling of the decay constants is found.

Uncertainties in the improvement and renormalisation coefficients are seen to have a large effect on the decay constants. The ALPHA and Bhattacharya *et al.* determinations of the improvement coefficients produce different values of f_K , especially at $\beta = 6.0$. However, the f_K do converge as the lattice spacing decreases. Computed quantities may well differ at $\mathcal{O}(a^2)$ at fixed lattice spacing but should agree in the continuum limit.

The heavy-light decay constants are found to be around 10% higher than, but compatible with, the world average [9] of quenched lattice determinations. HQS relations are found to be well satisfied by the data at both values of the coupling. The systematic uncertainties in the calculation have been discussed and, with the exception of quenching, estimated. The issue of $\mathcal{O}(a^2 m^2)$ effects in the heavy quark extrapolations has been considered in various ways. In particular, the difference of the decay constants between $\beta = 6.0$ and $\beta = 6.2$ at fixed heavy quark mass, as a measure of discretisation effects, is small and remains so during the heavy quark extrapolation. The static point value of the decay constant, despite the lack of improvement, suggests the extrapolation is under control. Furthermore, HQS relations are satisfied and there is agreement between f_{B_s}/f_B computed by extrapolating the ratio and by taking the ratio of the extrapolated decay constants.

The quenched spectrum of D mesons has been determined, shows good scaling, and is found to be in broad agreement with experimental data.

Acknowledgements

Support from EPSRC grant GR/K41663, and PPARC grants GR/L29927, GR/L56336 and PPA/G/S/1999/00022 is acknowledged. DGR acknowledges PPARC and the DOE (contract DE-AC05-84ER40150), and thanks FNAL for hospitality during part of this work. LDD thanks MURST for financial help. CMM acknowledges PPARC grant PPA/P/S/1998/00255. GNL thanks the Edinburgh Department of Physics and Astronomy for financial support. The authors thank Sara Collins, Christine Davies, Martin Kurth, Craig McNeile, Sinéad Ryan and Hartmut Wittig for useful discussions.

References

- [1] M. Neubert and B. Stech, in *Heavy Flavours II*, Vol. 15 of *Advanced Series on High Energy Physics*, edited by A.J. Buras and M. Lindner (World Scientific, Singapore, 1998), Chap. 4, pp. 294–344.
- [2] M. Ciuchini, R. Contino, E. Franco, and G. Martinelli, *Eur. Phys. J.* **C9**, 43 (1999).
- [3] C. Bernard, T. Draper, G. Hockney and A. Soni, *Phys. Rev. D* **38**, 33540 (1988).
- [4] M. Gavela *et al.*, *Phys. Lett. B* **206**, 113 (1988).
- [5] E. Eichten, *Nucl. Phys. B (Proc. Suppl.)* **4**, 170 (1988).
- [6] G. Lepage and B. Thacker, *Nucl. Phys. B (Proc. Suppl.)* **4**, 199 (1987).
- [7] A.X. El-Khadra, A.S. Kronfeld, and P.B. Mackenzie, *Phys. Rev. D* **55**, 3933 (1997).
- [8] A.S. Kronfeld, *Phys. Rev. D* **62**, 014505 (2000).
- [9] C. Bernard, *Nucl. Phys. B (Proc. Suppl.)* **94**, 159 (2001).
- [10] B. Sheikholeslami and R. Wohlert, *Nucl. Phys. B* **259**, 572 (1985).
- [11] M. Kurth and R. Sommer, *Nucl. Phys. B* **597**, 488 (2001).
- [12] M. Creutz, *Phys. Rev. D* **36**, 2394 (1987).
- [13] F. Brown and T. Woch, *Phys. Rev. Lett.* **58**, 163 (1987).
- [14] N. Cabibbo and E. Marinari, *Phys. Lett. B* **119**, 387 (1982).
- [15] UKQCD Collaboration, K.C. Bowler *et al.*, *Phys. Rev. D* **62**, 054506 (2000).
- [16] B. Efron, *SIAM Review* **21**, 460 (1979).
- [17] M. Lüscher, S. Sint, R. Sommer and P. Weisz, *Nucl. Phys. B* **478**, 365 (1996).

- [18] M. Lüscher, S. Sint, R. Sommer, P. Weisz, H. Wittig and U. Wolff, Nucl. Phys. B (Proc. Suppl.) **53**, 905 (1997).
- [19] M. Lüscher, S. Sint, R. Sommer, and H. Wittig, Nucl. Phys. B **491**, 344 (1997).
- [20] R. Sommer, Nucl. Phys. Proc. Suppl **60A**, 279 (1998).
- [21] M. Guagnelli and R. Sommer, Nucl. Phys. B (Proc. Suppl.) **63A-C**, 886 (1998).
- [22] T. Bhattacharya *et al.*, Phys. Lett. B **461**, 79 (1999).
- [23] T. Bhattacharya, R. Gupta, W. Lee, S. Sharpe, Nucl. Phys. B (Proc. Suppl.) **83-84**, 851 (2000).
- [24] T. Bhattacharya, R. Gupta, W. Lee, S. Sharpe, Phys. Rev. D **62**, 074505 (2001).
- [25] S. Sint and P. Weisz, Nucl. Phys. B **502**, 251 (1997).
- [26] G. Lepage and P. Mackenzie, Phys. Rev. D **48**, 2250 (1993).
- [27] S. Collins and C.T.H Davies, Nucl. Phys. B (Proc. Suppl.) **94**, 608 (2001).
- [28] UKQCD Collaboration, K.C. Bowler *et al.*, Phys. Lett. B **486**, 111 (2000).
- [29] UKQCD Collaboration, C.M. Maynard, Nucl. Phys. B (Proc. Suppl.) **94**, 367 (2001).
- [30] UKQCD Collaboration, in preparation.
- [31] G.M. de Divitiis and R. Petronzio, Phys. Lett. B **419**, 311 (1998).
- [32] M. Kurth, private communication.
- [33] R. Sommer, Phys. Rept. **275**, 1 (1996).
- [34] D.E. Groom *et al.*, Eur. Phys. J. C **15**, 1 (2000).
- [35] P. Boyle, `hep-lat/9903033`.
- [36] P. Lacock, A. McKerrell, C. Michael, I.M. Stopher and P.W. Stephenson, Phys. Rev. D **51**, 6403 (1995).
- [37] R. Sommer, Nucl. Phys. **B411**, 839 (1994).
- [38] M. Guagnelli, R. Sommer and H. Wittig, Nucl. Phys. B **535**, 389 (1998).
- [39] M. Neubert, Phys. Rev. D **45**, 2451 (1992).
- [40] M. Neubert, Phys. Rev. D **46**, 1076 (1992).
- [41] S. Bethke, Nucl. Phys. B (Proc. Suppl.) **54A**, 314 (1997).
- [42] J. Garden, J. Heitger, R. Sommer and H. Wittig, Nucl. Phys. B **571**, 237 (2000).
- [43] S. Capitani, M. Lüscher, R. Sommer, H. Wittig, Nucl. Phys. B **544**, 669 (1999).
- [44] C. Bernard, J. Labrenz and A. Soni, Phys. Rev. D **49**, 2536 (1994).

- [45] G.P. Lepage, Nucl. Phys. B (Proc. Suppl.) **26**, 45 (1992).
- [46] A. Kronfeld, Nucl. Phys. B (Proc. Suppl.) **30**, 445 (1993).
- [47] A.X. El-Khadra *et al.*, Phys. Rev. D **58**, 014506 (1998).
- [48] C. Bernard *et al.*, Phys. Rev. Lett. **81**, 4812 (1998).
- [49] C. Bernard *et al.*, Nucl. Phys. B (Proc. Suppl.) **94**, 346 (2001).
- [50] A. Aoki *et al.*, Phys. Rev. Lett. **80**, 5711 (1998).
- [51] UKQCD Collaboration, L. Lellouch and C.J.D. Lin, [hep-ph/0011086](#).
- [52] D. Becirevic *et al.*, Phys. Rev. D **60**, 074501 (1999).
- [53] CP-PACS Collaboration, A. Ali Khan *et al.*, [hep-lat/0010009](#).
- [54] F. Muheim, private communication.